

# Inverse scattering with non-over-determined data

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## Abstract

The results of the author's theory of the inverse scattering with non-over-determined data are described.

## 1 Introduction

There is a large literature on inverse scattering, see [1] and references therein. We consider the potential scattering and the obstacle scattering.

The potential scattering problem consists of finding the scattering solution  $u(x, \alpha, k)$ :

$$[\nabla^2 + k^2 - q(x)]u = 0 \quad \text{in } \mathbb{R}^3, \quad (1)$$

$$u = u_0 + v, \quad u_0 = e^{ik\alpha \cdot x} \quad (2)$$

$$v_r - ikv = O(r^{-2}), \quad r \rightarrow \infty. \quad (3)$$

Here  $r := |x|$ ,  $\alpha \in S^2$ ,  $S^2$  is the unit sphere in  $\mathbb{R}^3$ ,  $q = q(x) \in L_{loc}^2(\mathbb{R}^3)$  is assumed to be real-valued and compactly supported. One has

$$v(x, \alpha, k) = A(\beta, \alpha, k) \frac{e^{ikr}}{r} + O(r^{-2}), \quad r \rightarrow \infty, \quad \beta = x/r. \quad (4)$$

The  $A(\beta, \alpha, k)$  is called the scattering amplitude,  $\beta \in S^2$  is the direction of the scattered wave.

The inverse scattering problem consists of finding  $q(x)$  from the scattering amplitude  $A$ . The function  $A$  is a function of five variables. It is easy to prove that this function known for all  $\alpha, \beta \in S^2$  and  $\forall k > 0$  determines  $q$  uniquely. In 1987 the author proved that a compactly supported potential  $q$  is uniquely determined by the fixed-energy scattering amplitude. More precisely, the values of  $A(\beta, \alpha, k_0)$  for  $\beta$  and  $\alpha$  running through fixed open subsets of  $S^2$  and  $k = k_0 > 0$  fixed determine a compactly supported  $q$  uniquely, see [3], [4], [5], [1]. The author also gave stability estimates for  $q$  in terms of the scattering amplitude, see [6], [1] and references therein.

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MSC: 35P25, 45Q0581U40.

Key words: inverse scattering with non-over-determined data.



However, the fixed-energy data is a function of four variable, while the  $q(x)$  is a function of three variables. The non-over-determined data are the values of the scattering amplitude which form a three-dimensional set. For example, the values  $A(-\alpha, \alpha, k)$  for all  $\alpha \in S^2$  and all  $k > 0$  is such a set. These are the back-scattering data at all energies. For compactly supported potentials the author proved uniqueness of the solution to the inverse scattering problem with the non-over-determined data  $A(-\alpha, \alpha, k)$  known for all  $k$  in an arbitrary small open subset of  $[0, \infty)$  and all  $\alpha$  in an arbitrary small open subset of  $S^2$ . The author proved that for a compactly supported potential these data determine uniquely the values of  $A(-\alpha, \alpha, k)$  for all  $k > 0$  and all  $\alpha \in S^2$ .

The other practically interesting example of non-over-determined data for which the author proved the uniqueness of the solution to the inverse scattering problem are the values of  $A(\beta, \alpha_0, k)$  known for all  $k$  in an arbitrary small open subset of  $[0, \infty)$  and all  $\beta$  in an arbitrary small open subset of  $S^2$ ,  $\alpha = \alpha_0$  being fixed.

These results are first published in [13], [14], [15] and in the monograph [1].

The obstacle scattering problem consists of finding the scattering solution  $u(x, \alpha, k)$ . Let  $D \subset \mathbb{R}^3$  be a bounded domain with a smooth connected boundary  $S$ ,  $D' := \mathbb{R}^3 \setminus D$ . Then

$$(\nabla^2 + k^2)u = 0 \quad \text{in } D', \quad u|_S = 0, \quad (5)$$

$$u = u_0 + v, \quad u_0 = e^{ik\alpha \cdot x} \quad (6)$$

$$v_r - ikv = O(r^{-2}), \quad r \rightarrow \infty. \quad (7)$$

One has

$$v(x, \alpha, k) = A(\beta, \alpha, k) \frac{e^{ikr}}{r} + O(r^{-2}), \quad r \rightarrow \infty, \quad \beta = x/r. \quad (8)$$

The non-over-determined data are the values of  $A(\beta, \alpha, k)$  on a two-dimensional subset of the set  $S^2 \times S^2 \times [0, \infty)$ . For example, such is the set  $\forall \beta \in S^2$ , a fixed  $\alpha = \alpha_0$  and a fixed  $k = k_0 > 0$ .

*The author proved that these non-over-determined data determine uniquely the surface  $S$  and the boundary condition on  $S$ .*

The boundary condition is assumed of the Dirichle, or Neumann, or impedance type. The impedance boundary condition is

$$u_N = \zeta u \quad \text{on } S. \quad (9)$$

Here  $\zeta = \zeta(s)$  is the boundary impedance and it is assumed that

$$\text{Im}\zeta \leq 0. \quad (10)$$

Assumption (10) guarantees uniqueness of the solution to the obstacle scattering problem, [11].

The uniqueness theorems for inverse obstacle scattering with non-over-determined data is proved by the author in [8], [1], [16].

*We now sketch this proof.*

Let us assume that two obstacles  $D_1$  and  $D_2$  generate the same scattering amplitude for all  $\beta \in S^2$ , a fixed  $\alpha$  and a fixed  $k = k_0 > 0$ , and prove that then  $D_1 = D_2$  and the boundary conditions are the same. If  $D_1 = D_2 := D$  then  $u_1 = u_2$  in  $D'$ , so  $u_1 = u_2$  and  $U_{1N} = u_{2N}$  on  $S := \partial D$ . Consequently, the boundary conditions are the same.

Let us prove that  $S_1 = S_2$  if  $A_1(\beta) = A_2(\beta)$  for all  $\beta \in S^2$ , where  $A_j(\beta) := A_j(\beta, \alpha_0, k_0)$ ,  $j = 1, 2$ .

If  $A_1(\beta) = A_2(\beta)$  then  $u_1(x, \alpha_0, k_0) = u_2(x, \alpha_0, k_0)$  for all  $x \in D'_{12} := \mathbb{R}^3 \setminus (D_1 \cup D_2)$ . This follows from Lemma 1.2.15 in [1], p.47. Let  $D^{12} := D_1 \cap D_2$ ,  $S_{12} := \partial D_{12}$ ,  $S^{12} := \partial D^{12}$ . One has  $u_1 = u_2 := u$  in  $\mathbb{R}^3 \setminus D^{12}$ . By Green's formula one gets

$$u = u_0 - \int_{S_1} g(x, s) u_N ds, \quad x \in D'_1 \quad (11)$$

and

$$u = u_0 - \int_{S_2} g(x, s) u_N ds, \quad x \in D'_2. \quad (12)$$

Since  $u$  and  $u_0$  are defined in  $\mathbb{R}^3 \setminus D^{12}$ , the integrals in (11) and (12) are also defined in  $\mathbb{R}^3 \setminus D^{12}$ . Consequently one obtains

$$\int_{S_1} g(x, s) u_N ds = \int_{S_2} g(x, s) u_N ds \quad x \in D_{12} \setminus D^{12}. \quad (13)$$

By Green's formula one has

$$u(x) = \int_{S_2} g(x, s) u_N ds - \int_{S_1} g(x, s) u_N ds = 0, \quad x \in D_{12} \setminus D^{12}, \quad (14)$$

where formula (13) was used.

Since  $u$  is an analytic function of  $x$  in  $\mathbb{R}^3 \setminus D^{12}$  and vanishes in  $D_{12} \setminus D^{12}$  it must vanish everywhere in  $\mathbb{R}^3 \setminus D^{12}$ .

This is a contradiction since

$$\lim_{|x| \rightarrow \infty} |u(x, \alpha_0, k_0)| = 1.$$

This contradiction proves that  $D_1 = D_2$ , so  $S_1 = S_2$ . The proof is complete.  $\square$

A study of the inverse scattering problems with non-over-determined data is of principal interest because these are the minimal data from which the unknown scatterer can be uniquely determined.

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